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# Arnold's conjecture and symplectic reduction 

A. Ibort, C. Martínez Ontalba<br>Dpto. de Física Teórica., Univ. Complutense de Madrid, 28040 Madrid, Spain

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#### Abstract

Fortune (1985) proved Arnold's conjecture for complex projective spaces, by exploiting the fact that $\mathbb{C} \mathbb{P}^{n-1}$ is a symplectic quotient of $\mathbb{C}^{n}$. In this paper, we show that Fortune's approach is universal in the sense that it is possible to translate Arnold's conjecture on any closed symplectic manifold $(Q, \Omega)$ to a critical point problem with symmetry on loops in $\mathbb{R}^{2 n}$ with its standard symplectic structure.


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## 1. Introduction

Let ( $Q, \Omega$ ) be a closed symplectic manifold. A symplectic diffeomorphism $\varphi$ of ( $Q, \Omega$ ) is said to be exact if it can be obtained by integrating a time-dependent hamiltonian vector field. More precisely, $\varphi$ will be exact if there exists a smooth time-dependent hamiltonian $h: Q \times[0,1] \rightarrow \mathbb{R}$ such that, defining $\varphi_{t}$ by

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} \varphi_{t}=X_{h_{t}} \circ \varphi_{t}, \quad \varphi_{0}=\mathrm{id}_{Q} \tag{1}
\end{equation*}
$$

where $X_{h_{1}}$ satisfies $\mathrm{i}\left(X_{h_{t}}\right) \Omega=-\mathrm{d} h_{t}$, one has $\varphi=\varphi_{1}$.
In the seventies, Arnold [1,2] conjectured that any exact symplectomorphism $\varphi_{1}$ of a closed symplectic manifold ( $Q, \Omega$ ) must have at least as many fixed points as the minimal number of critical points for a smooth real-valued function on $Q$. Moreover, if all fixed points are nondegenerate, then the lower bound is given by the minimal number of critical points for a Morse function on $Q$.

The conjecture has been proved for several classes of manifolds, by using different methods and techniques. (A more complete overview of the literature can be found in [3].)

The most general results come from Floer's work [3]. Floer proved Arnold's conjecture for manifolds $(Q, \Omega)$ for which the class $[\Omega] \in H^{2}(Q, \mathbb{R})$ vanishes on $\pi_{2}(Q)$, and the nondegenerate part of the conjecture for monotone manifolds, i.e., manifolds for which [ $\Omega$ ] is positively proportional to the first Chern class $c_{1}(Q)$ on $\pi_{2}(Q)$. A crucial ingredient in Floer's proof is the existence of an action functional or, failing that, of a well-defined vector field on the space of loops whose critical points correspond to closed orbits of the hamiltonian system on $Q$. Floer's techniques have been recently refined (see e.g. [9,13,8]) to extend his results to a wider range of manifolds.

A quite different approach is due to Fortune and Weinstein [5,4]. In the case of $\mathbb{C} \mathbf{P}^{n-1}$, the action functional is multiple valued. This difficulty disappears when one considers the hamiltonian system on $\mathbb{C} \mathbf{P}^{n-1}$ as the reduction, in the sense of Marsden and Weinstein, of a hamiltonian system on $\mathbb{C}^{n}$. The problem can be reduced to that of finding certain families of critical orbits of the restriction of an $S^{1}$-invariant action functional, defined in the loop space of $\mathbb{C}^{n}$, to a given invariant submanifold.

The same approach has been recently used by Oh [12] and Givental [6] to get estimates for the minimal number of fixed points of exact symplectomorphisms of $\mathbb{T}^{2 n} \times \mathbb{C P}^{k}$ and toric symplectic manifolds (i.e., symplectic quotients $\mathbb{C}^{n} / / \mathbb{T}^{k}$ with respect to certain linear torus actions), respectively.

In this paper, we will show how, making use of a suitable inverse reduction, Arnold's conjecture on any closed symplectic manifold ( $Q, \Omega$ ) can be formulated as a critical point problem with symmetry on loops in a canonical cotangent bundle ( $T^{*} P, \mathrm{~d} \theta_{P}$ ) or, even more, as a critical point problem with symmetry on loops in some $\mathbb{R}^{2 n}$ with its standard symplectic structure. The groups involved here are $\mathbb{T}^{k}$, for the problem in ( $T^{*} P, \mathrm{~d} \theta_{P}$ ), and a product $\mathbb{T}^{k} \times \mathbb{R}^{l}$, for the problem in $\mathbb{R}^{2 n}$. We will provide a detailed proof of these facts as well as a complete discussion of the resulting variational problems, extending previous results announced in [10].

The proof relies on the fact, showed in [7], that every symplectic manifold ( $Q, \Omega$ ), with $\Omega$ of finite integral rank, can be realized as a symplectic reduction (although not always as a Marsden-Weinstein reduction) of some $\mathbb{R}^{2 n}$ with its standard symplectic structure.

In Section 2 we will extend the use of inverse reduction in [4] to the more general context of the Marsden-Weinstein reduction of a symplectic manifold ( $M, \omega$ ) with respect to the action of a connected abelian Lie group $G$. We will see how to express the fixed point problem on the reduced manifold as a fixed point problem with symmetry on, the usually simpler one, $(M, \omega)$. Then, in Section 1, we will combine these results with Gotay and Tuynman's theorem and, with an appropriate lifting of the hamiltonian system on ( $Q, \Omega$ ), we will prove our main result (cf. Theorems 2 and 3 ).

## 2. Fixed points and inverse reduction

Let $G$ be a Lie group acting symplectically on a symplectic manifold ( $M, \omega$ ). Let us assume that the action admits an $\mathrm{Ad}^{*}$-equivariant momentum map $J: M \rightarrow \mathcal{G}^{*}$, where $\mathcal{G}^{*}$ denotes the dual of the Lie algebra $\mathcal{G}$ of $G$.

If $\mu \in \mathcal{G}^{*}$ and $J^{-1}(\mu)$ is a submanifold of $M$, then there is an induced action of the stabilizer group of $\mu, G_{\mu}$, on $J^{-1}(\mu)$. We will assume that the quotient space $M_{\mu}=J^{-1}(\mu) / G_{\mu}$ is a smooth manifold and the projection $\pi_{\mu}: J^{-1}(\mu) \rightarrow M_{\mu}$ is a smooth submersion. Under these hypotheses, there is an induced symplectic form $\omega_{\mu}$ on $M_{\mu}$ and ( $M_{\mu}, \omega_{\mu}$ ) is known as the Marsden-Weinstein reduction of $(M, \omega)$ relative to the group action.

In what follows, we will restrict ourselves to the action $\Phi: G \times M \rightarrow M$ of a finitedimensional connected abelian Lie group $G$ on $M$. (Notice that $G$ must be isomorphic to a product $\mathbb{T}^{k} \times \mathbb{R}^{l}$.)

Let us consider a given time-dependent hamiltonian $H_{t}$ on $M_{\mu}$ with associated timedependent hamiltonian vector field $X_{H_{t}}$. We are looking for closed integral curves of $X_{H_{l}}$, i.e., closed solutions of

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} u(t)=X_{H_{t}}(u(t)), \quad u(0)=m \tag{2}
\end{equation*}
$$

for any $m \in M_{\mu}$.
Now assume that there exists a time-dependent hamiltonian $\tilde{H}: M \times[0,1] \rightarrow \mathbb{R}$ on $M$ such that each $\tilde{H}_{t}$ is a $G$-invariant extension to $M$ of the pull-back $\pi_{\mu}^{*} H_{t}$ and let $X_{\tilde{H}_{t}}$ be the corresponding hamiltonian vector field. It is easily seen that $X_{\tilde{H}_{l}}$ is tangent to $J^{-1}(\mu)$ and it projects on $X_{H_{t}}$.

Consider any $x \in J^{-1}(\mu)$ and let $\tilde{\sigma}_{x}$ be the solution of the initial value problem in $M$ :

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} u(t)=X_{\tilde{H}_{t}}(u(t)), \quad u(0)=x . \tag{3}
\end{equation*}
$$

Then $\tilde{\sigma}_{x}$ lies in $J^{-1}(\mu)$ and $\sigma_{m}=\pi_{\mu} \circ \tilde{\sigma}_{x}$ is the solution of (2) with $m=\pi_{\mu}(x)$. For each $g \in G, \Phi_{g} \circ \tilde{\sigma}_{x}=\tilde{\sigma}_{\Phi_{g}(x)}$ will also be an integral curve of $X_{\tilde{H}_{t}}$ projecting on $\sigma_{m}$. In fact, there is a one-to-one correspondence between integral curves $\sigma_{m}$ of $X_{H_{l}}$ and families of integral curves of̂ $X_{\tilde{H}_{l}}$ with initial values at the points of the orbit $\pi_{\mu}^{-1}(m)$.

The curve $\sigma_{m}$ will be closed if and only if each $\tilde{\sigma}_{x}$ in the corresponding family satisfies $\tilde{\sigma}_{x}(0)=\Phi_{g_{0}}\left(\tilde{\sigma}_{x}(1)\right)$ for certain $g_{0} \in G$. Since $G$ is a connected abelian group, the exponential mapping $\exp : \mathcal{G} \rightarrow G$ is onto. If we pick any $\xi \in \exp ^{-1}\left(g_{0}\right)$ and define $\tilde{\sigma}_{x}^{\xi}(t)=\Phi_{g_{\xi}(t)}\left(\tilde{\sigma}_{x}(t)\right)$, where $g_{\xi}$ denotes the curve $t \mapsto \exp (t \xi)$ in $G$, then $\tilde{\sigma}_{x}^{\xi}$ is a closed integral curve of $X_{\tilde{H}_{t}}+\xi_{M}$, with $\xi_{M}$ being the infinitesimal generator of the action on $M$ corresponding to $\xi \in \mathcal{G}$, as we next prove.

It is obvious that $\tilde{\sigma}_{x}^{\xi}(0)=\tilde{\sigma}_{x}^{\xi}(1)=x$.
Now, differentiation of the expression of $\tilde{\sigma}_{x}^{\xi}(t)$ with respect to $t$ yields

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} \tilde{\sigma}_{x}^{\xi}(t)=T_{\tilde{\sigma}_{x}(t)} \Phi_{g_{\xi}(t)}\left(X_{\tilde{H}_{t}}\left(\tilde{\sigma}_{x}(t)\right)\right)+T_{g_{\xi}(t)} \Phi_{\tilde{\sigma}_{x}(t)}\left(\dot{g}_{\xi}(t)\right) \tag{4}
\end{equation*}
$$

Since $X_{\tilde{H}_{t}}$ is $G$-invariant, the first term turns out to be $X_{\hat{H}_{t}}\left(\tilde{\sigma}_{x}^{\xi}(t)\right)$.
On the other hand, denoting by $L$ left translation in $G$ and by $e$ the unit element of $G$, we have

$$
\begin{align*}
T_{g_{\xi}(t)} \Phi_{\tilde{\sigma}_{x}(t)}\left(\dot{g}_{\xi}(t)\right) & =T_{e}\left(\Phi_{\tilde{\sigma}_{x}(t)} \circ L_{g_{\xi}(t)}\right)(\xi)=T_{e}\left(\Phi_{g_{\xi}(t)} \circ \Phi_{\tilde{\sigma}_{x}(t)}\right)(\xi) \\
& =T_{\tilde{\sigma}_{x}(t)} \Phi_{g_{\xi}(t)}\left(\xi_{M}\left(\tilde{\sigma}_{x}(t)\right)\right)=\xi_{M}\left(\tilde{\sigma}_{x}^{\xi}(t)\right) \tag{5}
\end{align*}
$$

and the desired result follows.

Thus, if $J_{\xi}=\langle J, \xi\rangle$ is the hamiltonian associated to $\xi_{M}$, then $\tilde{\sigma}_{x}^{\xi}$ is the solution of the initial value problem in $M$ :

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} u(t)=X_{\tilde{H}_{t}+J_{\xi}}(u(t)), \quad u(0)=x \tag{6}
\end{equation*}
$$

and it is closed.
Notice that $J_{\xi}$ is also $G$-invariant and the reduced vector field of $X_{\tilde{H}_{l}+J_{\xi}}$ is again $X_{H_{t}}$.
For a nonfree action ${ }^{1}$ the element $g_{0}$ mentioned above is not unique. It can be replaced by any $g \in g_{0} G_{x}$, where $G_{x}$ denotes the stabilizer group of $x$. Moreover, the correspondence $\xi \mapsto \tilde{\sigma}_{x}^{\xi}$ is not one-to-one. More explicitly, since the curves $\tilde{\sigma}_{x}^{\xi}$ satisfy $\Phi_{g} \circ \tilde{\sigma}_{x}^{\xi}=\tilde{\sigma}_{\Phi_{g}(x)}^{\xi}$, the stabilizer of $\tilde{\sigma}_{x}^{\xi}(t)$ is $G_{x}$ for each $t$. On the other hand,

$$
\begin{equation*}
\tilde{\sigma}_{x}^{\eta}(t)=\Phi_{g_{\eta-\xi}(t)}\left(\tilde{\sigma}_{x}^{\xi}(t)\right) \tag{7}
\end{equation*}
$$

and

$$
\begin{equation*}
\tilde{\sigma}_{x}^{\eta}=\tilde{\sigma}_{x}^{\xi} \Leftrightarrow \exp (t(\eta-\xi)) \in G_{x}, \quad \forall t \Leftrightarrow \eta-\xi \in \mathcal{G}_{x}, \tag{8}
\end{equation*}
$$

where $\mathcal{G}_{x}$ denotes the Lic algebra of $\boldsymbol{G}_{x}$, which can also be characterized as $\mathcal{G}_{x}=\{\xi \in$ $\left.\mathcal{G} \mid \xi_{M}(x)=0\right\}$.

Now we are ready to state the following proposition.
Proposition 1. To each fixed point $m$ of the exact symplectomorphism $\psi_{1}$ induced by the hamiltonian $H_{l}$ on $M_{\mu}$ there corresponds a family of closed curves in $J^{-1}(\mu): \mathcal{F}_{m, g_{0}}=$ $\left\{\tilde{\sigma}_{x}^{\xi} \in C^{\infty}\left(S^{1}, M\right) \mid \tilde{\sigma}_{x}^{\xi}\right.$ solves $\left.(6), x \in \pi_{\mu}^{-1}(m), \xi \in \exp ^{-1}\left(g_{0} G_{x}\right) \bmod \mathcal{G}_{x}\right\}$, for certain $g_{0} \in G$. This family is diffeomorphic to the product of an orbit $G / G_{x}$ and the projection of $\exp ^{-1}\left(G_{x}\right)$ on $\mathcal{G} / \mathcal{G}_{x}$.

## 3. Lifting to $\mathbb{R}^{2 n}$

Let ( $Q, \Omega$ ) be a closed symplectic manifold and consider a time-dependent hamiltonian $h_{t}$ on $Q$ with associated hamiltonian vector field $X_{h_{t}}$. Denoting by $\varphi_{t}$ the flow of $X_{h_{t}}$, we are concerned with the number of fixed points of $\varphi_{1}$.

As mentioned above, our purpose is to translate the fixed point problem on ( $Q, \Omega$ ) to a critical point problem on loops in $\mathbb{R}^{2 n}$. We have been motivated by the following theorem.

Theorem 1 (Gotay and Tuynman [7]). Every symplectic manifold ( $Q, \Omega$ ), with $\Omega$ of finite integral rank, can be realized as a reduction of some $\mathbb{R}^{2 n}$ with its standard symplectic structure.

Since $Q$ is compact, the condition of $\Omega$ having finite integral rank is automatically satisfied in our case.

[^0]On the other hand, reduction in Theorem 1 must be understood in the following sense.
If ( $M, \omega$ ) is a symplectic manifold and $N$ is a submanifold such that the pull-back $\omega_{N}$ of $\omega$ to $N$ has constant rank and ker $\omega_{N}$ is fibrating, then the quotient (symplectic) manifold $M_{N}=N / \operatorname{ker} \omega_{N}$ is called the reduction of $M$ by $N$.

Therefore, we cannot directly apply the results of Section 2 in order to prove our statement.
We will develop our proof in three stages, according to the scheme of proof of Theorem 1.
Step 1: Let us consider the cotangent bundle $\tau_{Q}: T^{*} Q \rightarrow Q$ and let $\theta_{Q}$ be the Liouville 1 -form on $T^{*} Q$. The zero section $Z_{Q}$ is a symplectic submanifold of ( $T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega$ ), and it is canonically symplectomorphic to ( $Q, \Omega$ ). Therefore, $(Q, \Omega)$ can be realized as the reduction of $\left(T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega\right)$ by $Z_{Q}$. This is the first step in the proof of Theorem 1.

Now, we need to lift the fixed point problem from $(Q, \Omega)$ to ( $T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega$ ). The next four lemmas will make the job.

Lemma 1. There is a one-to-one correspondence between symplectomorphisms of ( $Q, \Omega$ ) and symplectomorphisms of $\left(T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega\right)$ preserving $\theta_{Q}$.

Proof. It is well known that a diffeomorphism of $T^{*} Q$ is the lift of a diffeomorphism of $Q$ if and only if it preserves $\theta_{Q}$, the latter being called homogeneous diffeomorphisms of $T^{*} Q$.

Thus, to each symplectomorphism $\varphi$ of ( $Q, \Omega$ ) one can associate the homogeneous diffeomorphism

$$
\begin{equation*}
T^{*} \varphi^{-1}: T^{*} Q \rightarrow T^{*} Q, \quad\left(q, \beta_{q}\right) \mapsto\left(\varphi(q), \beta_{q} \circ T_{\varphi(q)} \varphi^{-1}\right) \tag{9}
\end{equation*}
$$

To show that this diffeomorphism preserves the whole symplectic form on $T^{*} Q$, it is enough to check that $\left(T^{*} \varphi^{-1}\right)^{*} \tau_{Q}^{*} \Omega=\tau_{Q}^{*} \Omega$. But this is clear from the property $\tau_{Q} \circ$ $T^{*} \varphi^{-1}=\varphi \circ \tau_{Q}$ and the fact that $\varphi$ is symplectic.

Conversely, given a homogeneous symplectomorphism $\psi$ of ( $T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega$ ), there exists a unique diffeomorphism $\varphi$ of $Q$ such that $\psi=T^{*} \varphi^{-1}$, and this diffeomorphism is symplectic, since: $\psi^{*}\left(\mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega\right)=\mathrm{d} \theta_{Q}+\tau_{Q}^{*} \Omega$ implies $\psi^{*} \tau_{Q}^{*} \Omega=\tau_{Q}^{*} \Omega$, that is, $\tau_{Q}^{*}\left(\varphi^{*} \Omega-\Omega\right)=0$. But $\tau_{Q}^{*}$ is injective, so that $\varphi^{*} \Omega=\Omega$, as was to be proved.

Lemma 2. There is a one-to-one correspondence between smooth hamiltonian isotopies $\varphi_{t}$ of $(Q, \Omega)$ and smooth hamiltonian isotopies $\psi_{t}$ of $\left(T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega\right)$ such that $\psi_{t}$ is homogeneous for each $t$.

Proof. Let $\varphi_{t}$ be the flow induced by a hamiltonian $h_{t}$ on ( $Q, \Omega$ ), and define $\psi_{t}=T^{*} \varphi_{t}^{-1}$. Then, the $\psi_{t}$ constitute a smooth family of homogeneous symplectomorphisms connecting $\psi_{1}$ with the identity map.

Now, let us consider the vector field defined by differentiating $\psi_{t}$ in $t$ :

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} \psi_{t}=\tilde{X}_{t} \circ \psi_{t} \tag{10}
\end{equation*}
$$

From $\tau_{Q} \circ \psi_{t}=\varphi_{t} \circ \tau_{Q}$, it follows that

$$
\begin{equation*}
T \tau_{Q}\left(\tilde{X}_{t}\right)=X_{h_{t}} \circ \tau_{Q} \tag{11}
\end{equation*}
$$

Therefore, $\mathrm{i}\left(\tilde{X}_{t}\right)\left(\tau_{Q}^{*} \Omega\right)=-\mathrm{d}\left(h_{t} \circ \tau_{Q}\right)$.
On the other hand,

$$
\begin{equation*}
\mathrm{i}\left(\tilde{X}_{t}\right) \mathrm{d} \theta_{Q}=\mathcal{L}_{\tilde{X}_{t}} \theta_{Q}-\mathrm{d}\left(\theta_{Q}\left(\tilde{X}_{t}\right)\right) \tag{12}
\end{equation*}
$$

where $\mathcal{L}$ stands for the Lie derivative. Since the $\psi_{t}$ are homogeneous, $\mathcal{L}_{\tilde{X}_{t}} \theta_{Q}=0$ and, finally,

$$
\begin{equation*}
\mathrm{i}\left(\tilde{X}_{t}\right)\left(\mathrm{d} \theta_{Q}+\tau_{Q}^{*} \Omega\right)=-\mathrm{d}\left(\theta_{Q}\left(\tilde{X}_{t}\right)+h_{t} \circ \tau_{Q}\right) \tag{13}
\end{equation*}
$$

Thus, the family $\psi_{t}$ is generated by the hamiltonian

$$
\begin{equation*}
H_{t}=h_{t} \circ \tau_{Q}+\theta_{Q}\left(\tilde{X}_{t}\right) . \tag{14}
\end{equation*}
$$

In terms of the original hamiltonian $h_{t}$

$$
\begin{equation*}
H_{t}\left(q, \beta_{q}\right)=h_{t}(q)+\beta_{q}\left(X_{h_{i}}(q)\right), \quad \forall\left(q, \beta_{q}\right) \in T^{*} Q \tag{15}
\end{equation*}
$$

Conversely, let $\psi_{t}$ be the flow induced by a hamiltonian $H_{t}$ on $\left(T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega\right)$ and suppose that, for each $t, \psi_{t}$ is homogeneous. Then, there exists a family $\varphi_{t}$ of symplectomorphisms of $(Q, \Omega)$ such that $\psi_{t}=T^{*} \varphi_{t}^{-1}, \forall t$.

If $j: Z_{Q} \hookrightarrow T^{*} Q$ denotes the inclusion of the zero section, then $H_{t} \circ j$ is a hamiltonian on the symplectic submanifold $Z_{Q}$, which in turn defines a hamiltonian $h_{t}$ on $Q$ by

$$
\begin{equation*}
H_{t} \circ j=h_{t} \circ \tau_{Q} \circ j \tag{16}
\end{equation*}
$$

A straightforward calculation shows that this hamiltonian generates the family $\varphi_{t}$.
The fixed points $\left(q, \beta_{q}\right)$ of $\psi_{1}$ are characterized by the two conditions $\varphi_{1}(q)=q$ and $\beta_{q} \circ T_{q} \varphi_{1}^{-1}=\beta_{q}$. The last one is equivalent to the existence of an eigenvector of $T_{q} \varphi_{1}$ with eigenvalue equal to 1 , namely, the vector defined by $\mathrm{i}\left(v_{q}\right) \Omega_{q}=\beta_{q}$. Hence, the following lemma holds.

Lemma 3. To each fixed point $q$ of $\varphi_{1}$, one can univocally associate a subspace $V_{q} \subset$ $T_{q}^{*} Q$ of fixed points of $\psi_{1}$, whose dimension is equal to the multiplicity of $q: m(q)=$ $\operatorname{dim} \operatorname{ker}\left(T_{q} \varphi_{1}-\mathrm{id}_{T_{q} Q}\right)$.

Moreover, for the nondegenerate case, we have the following lemma.
Lemma 4. $\psi_{1}$ is nondegenerate if and only if $\varphi_{1}$ is nondegenerate.
Under these circumstances, there is a one-to-one correspondence between fixed points of both symplectomorphisms.

Proof. Suppose that $\varphi_{1}$ is nondegenerate and let ( $q, \beta_{q}$ ) be a fixed point of $\psi_{1}$. By the preceding remarks, it must be $\beta_{q}=0_{q}$. Let $w_{\left(q, 0_{q}\right)} \in T_{\left(q, 0_{q}\right)} T^{*} Q$ be an eigenvector of $T_{\left(q, 0_{q}\right)} \psi_{1}$ with eigenvalue equal to 1 . The nondegeneracy of $q$ implies $w_{\left(q, 0_{q}\right)} \in \operatorname{ker} T_{\left(q, 0_{q}\right)} \tau_{Q}$. As a consequence, the 1 -form $\left.\mathrm{i}\left(w_{\left(q, 0_{q}\right)}\right) \mathrm{d}_{\left(q, 0_{q}\right)} \theta_{Q}+\left(\tau_{Q}^{*} \Omega\right)_{\left(q, 0_{q}\right)}\right)$ is a 1-form on $T_{\left(q, 0_{q}\right)} Z_{Q}$.

Using the canonical identification between $Z_{Q}$ and $Q$, it is not difficult to check that this 1-form defines a $T_{q} \varphi_{1}$-invariant 1 -form on $T_{q} Q$. Since the last 1 -form must be zero, we conclude that $w_{\left(q, 0_{q}\right)}=0$.

Conversely, suppose that $\psi_{1}$ is nondegenerate and consider a fixed point $q$ of $\varphi_{1}$. Then, $\left(q, 0_{q}\right)$ is a nondegenerate fixed point of $\psi_{1}$. Using again the identification between $Z_{Q}$ and $Q$, it is easily seen that to each vector $v_{q} \in T_{q} Q$ one can univocally associate a vector $w_{\left(q, 0_{q}\right)} \in T_{\left(q, 0_{q}\right)} T^{*} Q$, tangent to $Z_{Q}$, which must be $T_{\left(q, 0_{q}\right)} \psi_{1}$-invariant if $v_{q}$ is $T_{q} \varphi_{1}$-invariant. Since ( $q, 0_{q}$ ) is nondegenerate, $w_{\left(q, 0_{q}\right)}$, and hence $v_{q}$, must be zero.

The fixed point problem in ( $Q, \Omega$ ) can be lifted in this way to an equivalent fixed point problem in $\left(T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega\right)$. We must search for the fixed points of the exact symplectomorphism $\psi_{1}$, induced by a hamiltonian $H_{t}$ of the form (15) on $T^{*} Q$. In the case of general exact symplectomorphisms, we must count, instead of single fixed points, whole subspaces of fixed points of $\psi_{1}$.

Step 2: The symplectic manifold $\left(T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega\right)$ is obtained as a reduction, in the sense of Marsden and Weinstein, of a canonical cotangent bundle ( $T^{*} P, \mathrm{~d} \theta_{P}$ ) as follows (see [7]).

Since $\Omega$ has finite integral rank, $\Omega=\mu_{1} c_{1}+\cdots+\mu_{k} c_{k}$, where $c_{1}, \ldots, c_{k}$ are integral closed 2 -forms and $\mu_{1}, \ldots, \mu_{k}$ are real constants. According to the prequantization procedure of Kostant and Souriau, for each $i=1, \ldots, k$ there exists a principal $S^{1}$-bundle $\pi_{i}: P_{i} \rightarrow Q$, with a connection $A_{i}$ whose curvature satisfies $\mathrm{d} A_{i}=\pi_{i}^{*} c_{i}$. The fiber product of the $P_{i}, \pi: P \rightarrow Q$, is a principal $\mathbb{T}^{k}$-bundle with connection $A=\left(A_{1}, \ldots, A_{k}\right)$ and curvature $\mathrm{d} A=\left(\pi^{*} c_{1}, \ldots, \pi^{*} c_{k}\right)$.

Let us denote by $\Phi$ the action of $\mathbb{T}^{k}$ on $P$, and by $\Phi^{T^{*}}$ the cotangent lifting of $\Phi$ to ( $T^{*} P, \mathrm{~d} \theta_{P}$ ). There is a canonically defined momentum map $J: T^{*} P \rightarrow \mathbb{R}^{k}$ given by

$$
\begin{equation*}
J\left(p, \gamma_{p}\right)=\left(T_{e} \Phi_{p}\right)^{*} \gamma_{p}=\gamma_{p} \cap T_{e} \Phi_{p} \tag{17}
\end{equation*}
$$

where $e=(1, \ldots, 1)$ is the unit element of $\mathbb{T}^{k}$.
The cotangent bundle reduction theorem (see [11]) provides a symplectic diffeomorphism $\nu_{A, \mu}$ between the symplectic manifolds ( $T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*}\left(\mu_{1} c_{1}+\cdots+\mu_{k} c_{k}\right)$ ) and $\left(J^{-1}(\mu) / \pi^{k},\left(\mathrm{~d} \theta_{P}\right)_{\mu}\right)$, where $\mu=\left(\mu_{1}, \ldots, \mu_{k}\right)$.

Our purpose is to find an invariant extension $\hat{H}_{t}$ of $H_{t} \circ \nu_{A, \mu} \circ \pi_{\mu}$ to $T^{*} P$, where $H_{t}$ is the hamiltonian (15).

The existence of such an extension is guaranteed by the fact that $\mathbb{T}^{k}$ is compact. It should also be clear from Section 2 that any invariant extension allows us to lift the problem to $T^{*} P$. However, we will next describe a natural construction of $\hat{H}_{t}$ in terms of the connection $A$.

Given the connection $A$ on $P$, one can construct a lift $\pi_{A}: T^{*} P \rightarrow T^{*} Q$ of $\pi$ to the cotangent bundles, as explained in [14].

For each $p \in P$ we have an exact sequence

$$
\begin{equation*}
0 \longrightarrow \mathbb{R}^{k} \xrightarrow{T_{e} \Phi_{p}} T_{p} P \xrightarrow{T_{p} \pi} T_{\pi(p)} Q \longrightarrow 0, \tag{18}
\end{equation*}
$$

The connection $A$ gives a splitting of (18), which in turn defines a linear map $\Gamma_{p}: T_{\pi(p)} Q \rightarrow$ $T_{p} P$ (the image of $\Gamma_{p}$ consisting of all horizontal vectors in $T_{p} P$ ).

We also have a splitting of the dual exact sequence

$$
\begin{equation*}
0 \longleftarrow \mathbb{R}^{k} \stackrel{\left(T_{e} \Phi_{p}\right)^{*}}{\longleftarrow} T_{p}^{*} P \stackrel{\left(T_{p} \pi\right)^{*}}{\longleftarrow} T_{\pi(p)}^{*} Q \longleftarrow 0 \tag{19}
\end{equation*}
$$

and $\pi_{A}$ is defined fiberwise by

$$
\begin{equation*}
\pi_{A} \mid T_{p}^{*} P=\Gamma_{p}^{*} . \tag{20}
\end{equation*}
$$

The map $\pi_{A}$ is a lift of $\pi$ in the sense that it makes the diagram

commutative. Moreover, $\pi_{A}$ is constant along the orbits of the lifted action $\Phi^{T^{*}}$.
From the explicit construction of $\nu_{A, \mu}$ in [11], it is easily seen that the diagram

$$
\begin{array}{rll}
J^{-1}(\mu) & \xrightarrow{\pi_{\mu}} & J^{-1}(\mu) / \mathbb{T}^{k} \\
i_{\mu} \downarrow & & \downarrow \nu_{A, \mu}  \tag{22}\\
T^{*} P & \xrightarrow{\pi_{A}} & T^{*} Q
\end{array}
$$

where $i_{\mu}$ is the inclusion of $J^{-1}(\mu)$, also commutes.
From the commutativity of this diagram, one can immediately derive the following lemma.

Lemma 5. The time-dependent hamiltonian $\hat{H}_{t}=H_{t} \circ \pi_{A}$ is $a \mathbb{J}^{k}$-invariant extension of $H_{t} \circ \nu_{A, \mu} \circ \pi_{\mu}$ to $T^{*} P$.

In terms of the hamiltonian $h_{t}$ on $Q$, we have

$$
\begin{equation*}
\hat{H}_{t}\left(p, \gamma_{p}\right)=h_{t}(\pi(p))+\gamma_{p}\left(\Gamma_{p}\left(X_{h_{t}}(\pi(p))\right)\right), \quad \forall\left(p, \gamma_{p}\right) \in T^{*} P . \tag{23}
\end{equation*}
$$

On the other hand, if $V_{q}$ is the subspace of Lemma 3, then the points in $J^{-1}(\mu)$ projecting onto $V_{q}$ constitute the $\mathbb{T}^{k}$-orbits of points in $\left\langle\mu, A_{p}\right\rangle+\left(T_{p} \pi\right)^{*}\left(V_{q}\right) \subset T_{p}^{*} P$, for any $p \in \pi^{-1}(q)$. Indeed, if $\left(p, \gamma_{p}\right)$ belongs to $J^{-1}(\mu)$, then $\gamma_{p}$ must be of the form

$$
\begin{equation*}
\gamma_{p}=\left\langle\mu, A_{p}\right\rangle+\beta_{\pi(p)} \circ T_{p} \pi, \tag{24}
\end{equation*}
$$

where $\beta_{\pi(p)} \in T_{\pi(p)}^{*} Q$. Besides,

$$
\begin{equation*}
\left(\nu_{A, \mu} \circ \pi_{\mu}\right)\left(p, \gamma_{p}\right)=\pi_{A}\left(p, \gamma_{p}\right)=\left(\pi(p), \gamma_{p} \circ \Gamma_{p}\right)=\left(\pi(p), \beta_{\pi(p)}\right) . \tag{25}
\end{equation*}
$$

Thus,

$$
\begin{equation*}
\left(v_{A, \mu} \circ \pi_{\mu}\right)^{-1}\left(V_{q}\right)=\bigcup_{p \in \pi^{-1}(q)}\left(\left\langle\mu, A_{p}\right\rangle+\left(T_{p} \pi\right)^{*}\left(V_{q}\right)\right) \tag{26}
\end{equation*}
$$

The invariance of $A$ and $T \pi$ imply that this union is in fact a union of orbits of the action $\Phi^{T^{*}}$ on $T^{*} P$.

From the definition of $\hat{H}_{t}$, it is clear that the integral curves of the associated vector field $X_{\hat{H}_{t}}$ project on horizontal lifts to $P$ of integral curves of $X_{h_{t}}$.

Now, let $\hat{\sigma}_{\left(p, \gamma_{p}\right)}$ and $\hat{\sigma}_{\left(p, \gamma_{p}^{\prime}\right)}$ be two integral curves of $X_{\hat{H}_{t}}$ with (different) initial values in $\left\langle\mu, A_{p}\right\rangle+\left(T_{p} \pi\right)^{*}\left(V_{q}\right)$ and corresponding to the same fixed point $q$ in $Q$. Let $g_{0}$ and $g_{0}^{\prime}$ be the elements of $\mathbb{T}^{k}$ relating their initial and final values. Then it is easy to check that since both $\tau_{P} \circ \hat{\sigma}_{\left(p, \gamma_{p}\right)}$ and $\tau_{P} \circ \hat{\sigma}_{\left(p, \gamma_{p}^{\prime}\right)}$ are equal to the horizontal lift to $P$ of the curve $\varphi_{t}(q), g_{0}$ and $g_{0}^{\prime}$ must coincide.

Combining the results of Steps 1 and 2, and using Proposition 1 particularized to the case of a free torus action, we can state the following theorem.

Theorem 2. Let us consider the initial value problem in $T^{*} P$ :

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} u(t)=X_{\hat{H}_{t}+J_{\lambda}}(u(t)), \quad u(0)=\left(p, \gamma_{p}\right) \tag{27}
\end{equation*}
$$

with $\lambda=\left(\lambda_{1}, \ldots, \lambda_{k}\right) \in \mathbb{R}^{k}$. Then, to each fixed point $q$ of $\varphi_{1}$ there corresponds a family of closed curves in $J^{-1}(\mu) \subset T^{*} P: \mathcal{F}_{q, a}=\left\{\hat{\sigma}_{\left(p, \gamma_{p}\right)}^{\lambda} \in C^{\infty}\left(S^{1}, T^{*} P\right) \mid \hat{\sigma}_{\left(p, \gamma_{p}\right)}^{\lambda}\right.$ solves (27), $\left.\lambda_{i}=a_{i}(\bmod 2 \pi), \forall i, \pi(p)=q, \gamma_{p} \in\left\langle\mu, A_{p}\right\rangle+\left(T_{p} \pi\right)^{*}\left(V_{q}\right)\right\}$, for certain $a=$ $\left(a_{1}, \ldots, a_{k}\right) \in \mathbb{R}^{k}$. This family is diffeomorphic to $V_{q} \times \mathbb{T}^{k} \times \mathbb{Z}^{k}$.

By Hamilton's principle, the closed integral curves of $X_{\hat{H}_{t}+J_{\lambda}}$ are exactly the critical points of the action functional

$$
\begin{equation*}
S_{\hat{H}, \lambda}(u)=\int_{u} \theta_{P}-\int_{0}^{1}\left(\hat{H}_{t}(u(t))+J_{\lambda}(u(t))\right) \mathrm{d} t \tag{28}
\end{equation*}
$$

defined on closed loops $u$ in $T^{*} P$.
Notice that this functional is $\mathbb{T}^{k}$-invariant, where the (free) action of $\mathbb{T}^{k}$ on the loop space of $T^{*} P$ is the obvious one: $\left(\Phi_{g}(u)\right)(t)=\Phi_{g}(u(t))$.

We are looking for critical loops satisfying

$$
\begin{equation*}
J(u(t))=\mu, \quad \forall t \in S^{1} \tag{29}
\end{equation*}
$$

Since the hamiltonians considered are $\mathbb{T}^{k}$-invariant, this pointwise condition is equivalent on critical loops to the constraint $\mathcal{J}(u)=\mu$, where

$$
\begin{equation*}
\mathcal{J}(u)=\int_{0}^{1} J(u(t)) \mathrm{d} t \tag{30}
\end{equation*}
$$

takes its values in $\mathbb{R}^{k}$. This map $\mathcal{J}$ can be seen as a momentum mapping for the action of $\mathbb{T}^{k}$ on the loop space of $T^{*} P$.

Following [4], one can apply now Lagrange's multipliers theorem and identify the critical points above with the critical points of the restriction $f$ of $S_{\hat{H} .0}$ to $\mathcal{J}^{-1}(\mu)$. Notice that because the action of $\mathbb{T}^{k}$ is free, the map $\mathcal{J}$ will be a submersion at every $u \in \mathcal{J}^{-1}(\mu)$.

Let us compute now the difference between critical values of $f$ corresponding to the same fixed point in $Q$.

If $u$ is a critical loop of $S_{\hat{H}, \lambda}$ belonging to the family associated to a closed curve $\varphi_{t}(q)$ in $Q$, then:

$$
\begin{align*}
f(u)= & \int_{0}^{1}\left(\theta_{P}\right)_{u(t)}(\dot{u}(t)) \mathrm{d} t-\int_{0}^{1} \hat{H}_{t}(u(t)) \mathrm{d} t \\
= & \int_{0}^{1}\left(\theta_{P}\right)_{u(t)}\left(\lambda_{T^{*} P}(u(t))\right) \mathrm{d} t+\int_{0}^{1} u(t)\left(T_{u(t)} \tau_{P}\left(X_{\hat{H}_{t}}(u(t))\right)\right) \mathrm{d} t \\
& -\int_{0}^{1} \hat{H}_{t}(u(t)) \mathrm{d} t=\int_{0}^{1} J_{\lambda}(u(t)) \mathrm{d} t-\int_{0}^{1} h_{t}\left(\left(\pi \circ \tau_{p} \circ u\right)(t)\right) \mathrm{d} t \\
= & \langle\mu, \lambda\rangle-\int_{0}^{1} h_{t}\left(\varphi_{t}(q)\right) \mathrm{d} t \tag{31}
\end{align*}
$$

and the difference between any two critical values of $f$ corresponding to a fixed point $q \in Q$ is of the form $\langle\mu, \eta\rangle$ with $\eta \in \exp ^{-1}(e)=\mathbb{Z}^{k}$.

Thus, to each fixed point of the hamiltonian system on $Q$ there corresponds a tower of critical values of $f$, parametrized by the set $\left\{\langle\mu, \eta\rangle \mid \eta \in \mathbb{Z}^{k}\right\}$.

Step 3: First of all, we will briefly recall the last step in Gotay and Tuynman's proof.
They first take a closed equivariant embedding $P \hookrightarrow \mathbb{R}^{n}$ in some orthogonal representation space for $\mathbb{T}^{k}$. (The existence of such an embedding is a well-known theorem of Mostow and Palais.)

In order to obtain ( $T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega$ ) as a Marsden-Weinstein reduction of the corresponding $T^{*} \mathbb{R}^{n} \simeq \mathbb{R}^{2 n}$, they first show that the ideal in $C^{\infty}\left(\mathbb{R}^{n}\right)$ of all functions vanishing on $P$ is generated over $C^{\infty}\left(\mathbb{R}^{n}\right)$ by a finite collection $\left\{f_{1}, \ldots, f_{l}\right\}$ of $\mathbb{T}^{k}$-invariant functions.

Then, they define, for each $i=1, \ldots, l$, a hamiltonian $F_{i}=f_{i} \circ \tau_{n}$, where $\tau_{n}: \mathbb{R}^{2 n} \rightarrow \mathbb{R}^{n}$ is the canonical cotangent projection. The flows of the corresponding hamiltonian vector fields define a symplectic $\mathbb{R}^{l}$-action with momentum map $F=\left(F_{1}, \ldots, F_{l}\right)$. The MarsdenWeinstein reduction of $\mathbb{R}^{2 n}$ relative to this action, $F^{-1}(0) / \mathbb{R}^{l}$, is symplectomorphic to ( $T^{*} P, \mathrm{~d} \theta_{P}$ ).

Since the functions $f_{i}$ have been chosen $\mathbb{T}^{k}$-invariant, the $\mathbb{R}^{l}$-action commutes with the cotangent lifting of the $\mathbb{T}^{k}$-action on $\mathbb{R}^{n}$, and there is a well-defined symplectic ( $\mathbb{T}^{k} \times \mathbb{R}^{l}$ )action with momentum map $K \times F$, where $K$ is the canonical momentum map for the torus action. Now, it follows that ( $T^{*} Q, \mathrm{~d} \theta_{Q}+\tau_{Q}^{*} \Omega$ ) is symplectomorphic to the reduced manifold $(K \times F)^{-1}(\mu ; \mathbf{0}) /\left(\mathbb{T}^{k} \times \mathbb{R}^{l}\right)$.

Notice that the $\mathbb{R}^{l}$-action is simply given by

$$
\begin{equation*}
\Psi_{\left(b_{1}, \ldots, b_{l}\right)}(x, y)=\left(x, y-\sum_{i=1}^{l} b_{i} d_{x} f_{i}\right), \quad \forall(x, y) \in \mathbb{R}^{2 n} \tag{32}
\end{equation*}
$$

and $F^{-1}(\mathbf{0})=T_{P}^{*} \mathbb{R}^{n}$. The projection $\pi_{P}: T_{P}^{*} \mathbb{R}^{n} \rightarrow T^{*} P$ is just the projection associated to the direct sum decomposition

$$
\begin{equation*}
T_{x}^{*} \mathbb{R}^{n}=T_{x}^{*} P \oplus \operatorname{span}\left[d_{x} f_{1}, \ldots, d_{x} f_{l}\right] \tag{33}
\end{equation*}
$$

for each $x \in P$, and it is $\mathbb{T}^{k}$-equivariant. The orbits of the $\mathbb{R}^{I}$-action on $T_{P}^{*} \mathbb{R}^{n}$ are diffeomorphic to $\mathbb{R}^{\text {codim } P}$.

On the other hand, $\pi_{P}\left(K^{-1}(\mu) \cap F^{-1}(\mathbf{0})\right)=J^{-1}(\mu)$ and the projection $(K \times F)^{-1}(\mu ; \mathbf{0})$ $\rightarrow(K \times F)^{-1}(\mu ; \boldsymbol{0}) /\left(\mathbb{T}^{k} \times \mathbb{R}^{l}\right)$ is $\left.\pi_{\mu} \circ \pi_{P}\right|_{K^{-1}(\mu) \cap F^{-1}(\mathbf{0})}$.

Next we prove the following lemma.
Lemma 6. There exists an $\left(\mathbb{T}^{k} \times \mathbb{R}^{l}\right)$-invariant extension $\tilde{H}_{t}$ of $H_{t} \circ v_{A}, \mu \circ \pi_{\mu} \circ$ $\left.\pi_{P}\right|_{K^{-1}(\mu) \cap F^{-1}(\mathbf{0})}$ to $\mathbb{R}^{2 n}$.

Proof. First notice that it will suffice to construct an $\mathbb{R}^{l}$-invariant extension $\tilde{H}_{l}$ of $\hat{H}_{l} \circ \pi_{P}$ which is also $\mathbb{T}^{k}$-invariant.

Since $\mathbb{T}^{k}$ is compact, we can always find a $\mathbb{T}^{k}$-invariant tubular neighborhood $U$ of $P$ in $\mathbb{R}^{n}$, a smooth $\mathbb{T}^{k}$-equivariant retraction $r: U \rightarrow P$, and a $\mathbb{T}^{k}$-invariant partition of unity $\left\{\rho_{0}, \rho_{1}\right\}$ subordinate to the open cover $\left\{U, \mathbb{R}^{n}-P\right\}$ of $\mathbb{R}^{n}$.

Now define

$$
\begin{equation*}
\tilde{H}: \mathbb{R}^{2 n} \times[0,1] \rightarrow \mathbb{R}, \quad((x, y), t) \mapsto \rho_{0}(x)\left(\hat{H}_{t} \circ \pi_{P}\right)(r(x), y) . \tag{34}
\end{equation*}
$$

It is easily seen that this hamiltonian gives an $\mathbb{R}^{l}$-invariant extension $\tilde{H}_{t}$ of $\hat{H}_{t} \circ \pi_{P}$, which is also $\mathbb{T}^{k}$-invariant.

As in Step 2, it is not hard to verify that the points in $\mathbb{R}^{2 n}$ projecting onto $V_{q}$ constitute the ( $T^{k} \times \mathbb{R}^{l}$ )-orbits of points belonging to $\left\langle\mu, A_{x}\right\rangle+\left(T_{x} \pi\right)^{*}\left(V_{q}\right)$, for any $x \in \pi^{-1}(q)$.

Now, we can state the following theorem.
Theorem 3. Let us consider the initial value problem in $\mathbb{R}^{2 n}$ :

$$
\begin{equation*}
\frac{\mathrm{d}}{\mathrm{~d} t} u(t)=X_{\tilde{H}_{t}+(K \times F)_{5}}(u(t)), \quad u(0)=(x, y) \tag{35}
\end{equation*}
$$

with $\xi=\left(\lambda_{1}, \ldots, \lambda_{k} ; b_{1}, \ldots, b_{l}\right) \in \mathbb{R}^{k} \times \mathbb{R}^{l}$. Then, to each fixed point $q$ of $\varphi_{1}$ there corresponds a family of closed curves in $(K \times F)^{-1}(\mu, \mathbf{0}) \subset \mathbb{R}^{2 n}: \tilde{\mathcal{F}}_{q, a}=$ $\left\{\tilde{\sigma}_{(x, y)}^{\xi} \in C^{\infty}\left(S^{1}, \mathbb{R}^{2 n}\right) \mid \tilde{\sigma}_{(x, y)}^{\xi}\right.$ solves $(35), \lambda_{i}=a_{i}(\bmod 2 \pi) \forall i, \pi(x)=q, y \in\left\langle\mu, A_{x}\right\rangle$ $\left.+\left(T_{x} \pi\right)^{*}\left(V_{q}\right)+\operatorname{span}\left\{d_{x} f_{1}, \ldots, d_{x} f_{l}\right\}\right\}$, for certain $a-\left(a_{1}, \ldots, a_{k}\right) \in \mathbb{R}^{k}$. This family is diffeomorphic to $V_{q} \times \mathbb{T}^{k} \times \mathbb{R}^{n-k-\operatorname{dim} Q} \times \mathbb{Z}^{k}$.

Remark. Notice that a closed curve $\tilde{\sigma}_{(x, y)}^{\xi}$ is determined by the data $(x, y)$ and $\lambda$.
As before, the fixed point problem in $Q$ can be translated to the problem of finding the critical points of the family of action functionals in the loop space of $\mathbb{R}^{2 n}$ :

$$
\begin{align*}
S_{\tilde{H}, \xi}(u) & =\int_{u} \theta_{0}-\int_{0}^{1}\left(\tilde{H}_{t}(u(t))+(K \times F)_{\xi}(u(t))\right) \mathrm{d} t \\
& =\int_{u} \theta_{0}-\int_{0}^{1} \tilde{H}_{t}(u(t)) \mathrm{d} t-\int_{0}^{1} K_{\lambda}(u(t)) \mathrm{d} t-\int_{0}^{1} F_{b}(u(t)) \mathrm{d} t \tag{36}
\end{align*}
$$

with $\xi=(\lambda ; b) \in \mathbb{R}^{k} \times \mathbb{R}^{l}$, subject to the constraints:

$$
\begin{equation*}
\mathcal{K}(u)=\int_{0}^{1} K(u(t)) \mathrm{d} t=\mu, \quad \mathcal{F}(u)=\int_{0}^{1} F(u(t)) \mathrm{d} t=0 . \tag{37}
\end{equation*}
$$

The action of $\mathbb{T}^{k} \times \mathbb{R}^{l}$ is not free now and $\mathcal{K}^{-1}(\mu) \cap \mathcal{F}^{-1}(\mathbf{0})$ will not be in general a submanifold. Nevertheless, one may still consider the components of $\lambda$ as Lagrange's multipliers, as we next explain.

Proposition 2. With the previous notation, the set $\mathcal{K}^{-1}(\mu)$ is a submanifold of codimension $k$ in the space of free loops on $\mathbb{R}^{2 n}$.

Proof. Let us consider a loop $u \in \mathcal{K}^{-1}(\mu)$. The tangent map $T_{u} \mathcal{K}$ fails to be surjective if and only if the stabilizer group of $u$ in $\mathbb{T}^{k}$ is not discrete, i.e. if and only if the corresponding Lie subalgebra $\left\{\lambda \in \mathbb{R}^{k} \mid \lambda_{\mathbb{R}^{2 n}} \circ u=0\right\}$ is not zero ( $\lambda_{\mathbb{R}^{2 n}}$ denotes, as is customary in this paper, the vector field associated to the element $\lambda$ in the Lie algebra of $\mathbb{T}^{k}$ by means of its action on $\mathbb{R}^{2 n}$ ).

If the stabilizer of $u$ is not discrete, it must contain a subgroup isomorphic to $S^{1}$. Hence, the stabilizer algebra must contain an element $\lambda \in(2 \pi \mathbb{Z})^{k}$ which generates the circle. This $\lambda$ will belong to the kernel of $\mu$, because

$$
\begin{equation*}
\lambda_{\mathbb{R}^{2 n}} \circ u=0 \Rightarrow K_{\lambda} \circ u=0 \Rightarrow \mathcal{K}_{\lambda}(u)=\langle\mu, \lambda\rangle=0 \tag{38}
\end{equation*}
$$

But $\mu=\left(\mu_{1}, \ldots, \mu_{k}\right)$ comes from the decomposition $\Omega=\mu_{1} c_{1}+\cdots+\mu_{k} c_{k}$, so that we may assume that $\mu_{1}, \ldots, \mu_{k}$ are independent over $\mathbb{Z}$, hence ker $\mu$ does not contain any nonzero $\lambda \in(2 \pi \mathbb{Z})^{k}$.

Therefore, $\mathcal{K}$ is a submersion at every $u \in \mathcal{K}^{-1}(\mu)$ and $\mathcal{K}^{-1}(\mu)$ is a submanifold.
Thus, the critical points of the family (36) satisfying (37) are exactly the critical points of the family of functionals $f_{b}, b \in \mathbb{R}^{l}$ defined by the restriction of the functionals

$$
\begin{equation*}
S_{\tilde{H}_{\cdot}(0, b)}=\int_{u} \theta_{0}-\int_{0}^{1} \tilde{H}_{t}(u(t)) \mathrm{d} t-\int_{0}^{1} F_{b}(u(t)) \mathrm{d} t \tag{39}
\end{equation*}
$$

to the submanifold $\mathcal{K}^{-1}(\mu)$, and satisfying $\mathcal{F}(u)=0$.
It must be noticed that the role played by the parameters $b$ is quite different to the role played by the Lagrange multipliers $\lambda$ because they do not produce any splitting of the critical
subsets. To each critical $\mathbb{T}^{k}$-orbit of the functional $f$ in Step 2 there corresponds a critical $\left(\mathbb{T}^{k} \times \mathbb{R}^{l}\right.$ )-orbit of a subfamily (parametrized by $\mathbb{R}^{l-\operatorname{codim} P}$ ) of functionals $f_{b}$.

Finally, a short computation shows that the critical values of the functionals $f_{b}$ corresponding to the same fixed point $q \in Q$ are again of the form

$$
\begin{equation*}
f_{b}(u)=\langle\mu, \lambda\rangle-\int_{0}^{1} h_{t}\left(\varphi_{t}(q)\right) \mathrm{d} t \tag{40}
\end{equation*}
$$

and hence they are arranged in a tower parametrized again by the set $\left\{\langle\mu, \eta\rangle \mid \eta \in \mathbb{Z}^{k}\right\}$.

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[^0]:    ${ }^{1}$ In Proposition 1 of [10], the action of $G$ on $J^{-1}(\mu)$ was assumed to be free. We derive here the corresponding result for a general action.

